

## Improvement of Flux Pinning Properties of YBCO Superconducting Films with BaMeO<sub>3</sub> (Me:Zr, Mn, Hf and Ir) Perovskite Nanoparticles by TFA-MOD Technique

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### Abstract

Superconducting materials increase performances of machines by intensively using in MRI in medicine, thermal bolometer, energy storage systems in transformer, magnetic separators, levitation, NMR, generators, engines, cables, superconducting wires and tapes, accelerators, electromagnets, and electronic transistors. In these applications, the high temperature superconductors (HTS) possess very high critical current density ( $J_c$ ) values at low temperature even in high magnetic fields. In these kinds of studies, many kinds of crystalline defects, such as fine precipitates of non-superconducting phases, dislocations, vacancies, grain boundaries, twin boundaries and so on, are considered to act as pinning centers. However, the  $J_c$  values rapidly decrease with increasing temperature in magnetic field. The main reason of the  $J_c$  depression is the lack of effective pinning in the structure. Depending on these reason, a novel technology has been developed by means of a nanostructure engineering to create artificial pinning centers in HTS materials. With this regard, increasing critical current density and improvement of flux pinning properties of YBa<sub>2</sub>Cu<sub>3</sub>O<sub>6.57</sub> (YBCO) superconducting films with BaMeO<sub>3</sub> (Me:Zr, Mn, Hf and Ir) perovskite nanodots, nanorods or nanoparticles, as pinning centers, on SrTiO<sub>3</sub> substrate were aimed.

### Key words

High temperature superconductors; YBCO; BaMeO<sub>3</sub>; Flux pinning; STO; TFA-MOD

## TFA-MOD Tekniğiyle BaMeO<sub>3</sub> (Me:Zr, Mn, Hf ve Ir) Perovskit Nanopartiküllü YBCO Süperiletken Filmlerin Akı İğnelenmesi Özelliklerin Geliştirilmesi

### Özet

Süperiletken malzemeler tıpta MRI cihazlarında, enerji depolama sistemlerinde, termal kameralarda, manyetik ayırıştırıcılarda, analiz cihazlarında (NMR), jeneratörlerde, motorlarda, kablolarda, süperiletken tel ve şeritlerde, hızlandırıcı mıknatıslarda ve transistörlerde, cihazların performanslarını arttırmak amacı ile kullanılmaktadır. Bu uygulamalarda, yüksek sıcaklık süperiletken malzemeleri (HTS) düşük sıcaklık ve yüksek manyetik alanlarda, çok yüksek kritik akım yoğunluğu ( $J_c$ ) değerlerine sahiptir. Bu çeşit çalışmalarda, süperiletken olmayan fazların nano boyutta çökeltileri; dislokasyonlar, boşluklar, tane sınırları, ikiz sınırları gibi çoğu kristal hatalar, akı iğnelenmesi merkezleri olarak düşünülmektedir. Ancak  $J_c$  değerleri manyetik alanda artan sıcaklıkla hızlı bir şekilde düşmektedir.  $J_c$ 'nin düşmesinin ana sebebi etkili iğnelenme merkezlerinin olmamasıdır. Bu nedenlere bağlı olarak yeni teknoloji HTS malzemelerde yapısal iğnelenme merkezleri nano yapılı mühendislik ile geliştirilmektedir. Bunu sağlamak amacıyla TFA-MOD metodu kullanılarak, yapısal iğnelenme merkezleri olarak BaMeO<sub>3</sub> (Me: Zr, Hf, Ir, vs.) perovskit yapılı nanonoktalar, nanoçubuklar veya nanopartiküller şeklinde yapıların YBa<sub>2</sub>Cu<sub>3</sub>O<sub>6.57</sub> (YBCO) süperiletken filmlerin içine ilave edilerek kritik akım yoğunluğunun artırılması ve akı iğnelenmesi özelliklerinin geliştirilmesi hedeflenmiştir.

### Anahtar kelimeler

Yüksek sıcaklık süperiletkenler; YBCO; BaMeO<sub>3</sub>; Akı iğnelenmesi; STO; TFA-MOD

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### 1. Introduction

Production of long-length and low cost YBCO films has been an essential aim in the field of superconducting processing. Besides, YBCO films need to possess a high critical current density ( $J_c$ )

under high magnetic fields for practical applications as power and high field applications (Gutierrez *et al.*, 2007; Puig *et al.*, 2005). Improvement of their performance in magnetic fields has become one of the interesting topics on

high temperature superconductors (HTS). However, the ability for HTS to carry currents is significantly reduced in the presence of magnetic fields. In order to counteract this effect, various methods to increase the HTS current carrying abilities in magnetic fields by flux pinning have been developed through the pinning of the quantized flux lines by nanoscale crystalline defects and impurities (Strickland *et al.*, 2008; Engel *et al.*, 2007).

Vacuum deposition techniques like pulsed-laser deposition (PLD) and metal organic chemical vapour deposition (MOCVD) have been used to produce these nanocomposite YBCO films. However high cost and small exposure area of these deposition techniques have hindered the large scale production (Liu *et al.* 2008). Because of this reason, the more practical and cost effective chemical solution deposition method (CSD) is preferred recently.

The way for flux pinning in the structure can be the use of nanoparticles-modified substrate surfaces, mixing rare earth doping in to the structure and doping with BaMeO<sub>3</sub> nanoparticles. Among these approaches, Ba(Zr, Mn, Hf and Ir)O<sub>3</sub> nanoparticles in YBCO thin films are one of the most popular ones which can prevent the vortex motion at high fields. The amount of dopant in the structure needs to be high enough to generate the density of defects, which is needed to enhance flux pinning in magnetic fields. However, excess amounts of this nonsuperconducting content can suppress the self-field and in-field  $J_c$  values significantly (Chen *et al.*, 2009).

In the present work, undoped and BaMeO<sub>3</sub> doped YBCO thin films were grown with a non-vacuum, cost effective chemical solution deposition method. Their microstructures and dependence of the critical current on applied magnetic fields were investigated as a function of dopant concentration.

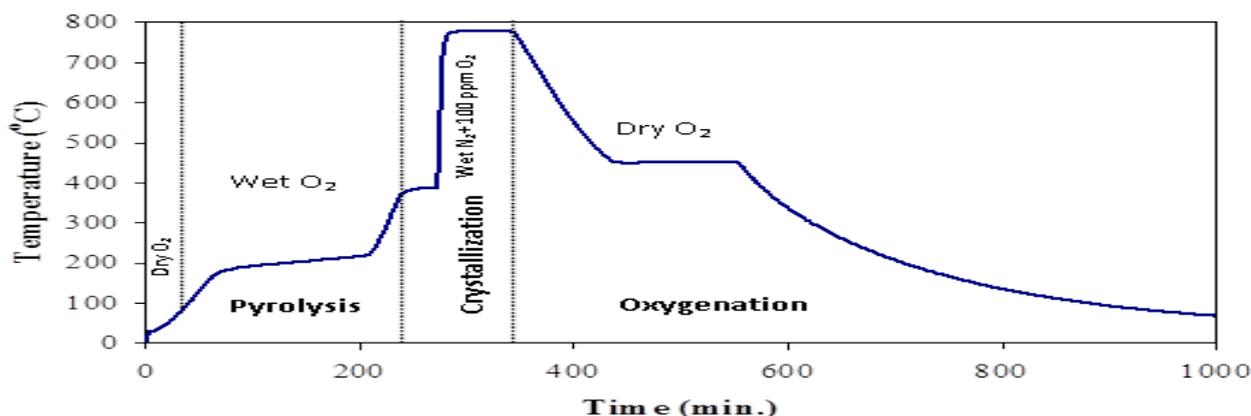
## 2. Material and Method

The precursor solutions of undoped and doped YBCO were prepared by dissolving acetates of Y, Ba and Cu with different amounts of Me(IV)-2,4-

pentanedionates in propionic acid with a quantity of TFA. It is followed by a solution purification process. The concentration of the solution was fixed to 0.25 M with addition of acetone and propionic acid with a certain ratio. Three different doped solutions were prepared with 6, 12 and 18 mol% BaMeO<sub>3</sub>. Film deposition on 10 mm x 10 mm STO (100) single crystal substrate was performed by spin coating technique at a rotation speed of 6000 rpm and acceleration speed of 6000 rpm s<sup>-1</sup>.

Coated samples were heat treated according to the profile shown in Figure 1. During the heat treatment process, dry gas treatment up to 60°C prevents the gel film from absorbing humidity that would deteriorate film integrity. Above 60°C, humidified O<sub>2</sub> is introduced to suppress the sublimation of Cu trifluoroacetate. Metal trifluoroacetates decompose and harmful gaseous residues are removed during the pyrolysis to give a precursor film. After that, the fluoride containing precursor film is fired at a maximum temperature of 780°C under humidified N<sub>2</sub> mixed with 100 ppm O<sub>2</sub> and then oxygenated at 450°C to obtain the YBCO superconducting film.

X-Ray Diffraction (XRD) was carried out using a Philips diffractometer with Co K $\alpha$  radiation to ascertain the phase purity of the undoped and doped films. Scanning Electron Microscopy (SEM, Philips XL20) was used to characterize the surface morphology of the final films. The critical transition temperature ( $T_c$ ) and critical current density ( $J_c$ ) of the films were measured by an inductive method. Transport measurements up to 6 T at 77 K on bridges of 0.8 mm length and 50  $\mu$ m widths were carried out with a physical properties measurement system (PPMS)

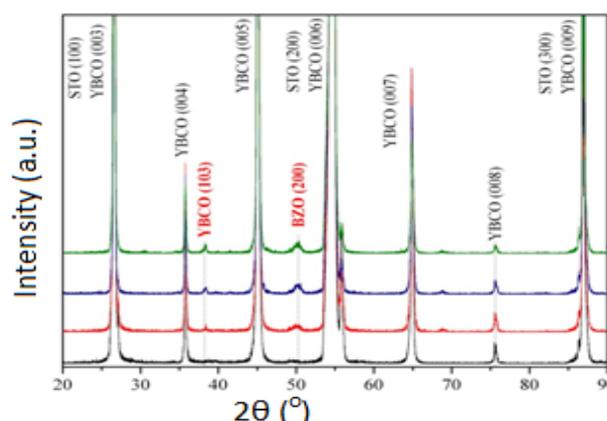


**Figure 1.** Heat treatment profile for the production of undoped and BaMeO<sub>3</sub> doped YBCO films. All films are crystallized at 780°C and oxygenated at 450°C.

### 3. Results and Discussion

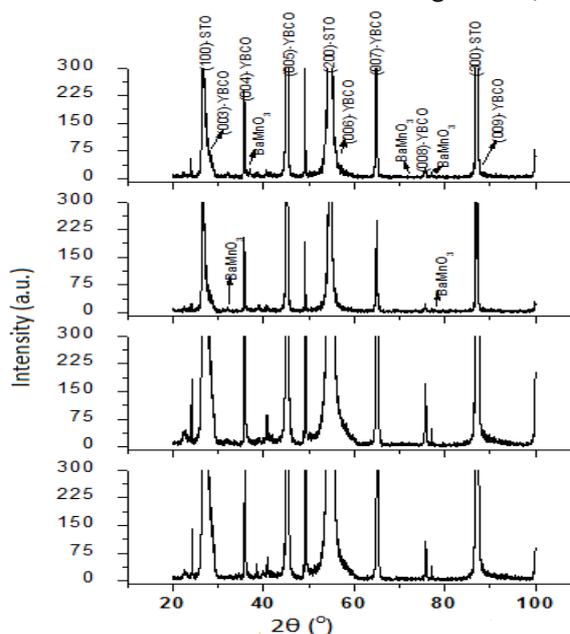
Figures 2, 3, 4 and 5 show the  $\theta$ - $2\theta$  diffraction patterns of undoped and BaMeO<sub>3</sub> doped YBCO films. In these patterns, the major peaks correspond to the (001) reflections of the YBCO phase and STO substrate which indicates that the YBCO film has a strong c-axis texture. Depending on the amount of dopant concentration, there is a slight increase at the BaZrO<sub>3</sub> (200) peak intensity. The structural integrity of the YBCO film can be easily affected by the doping process. In all doped samples (103) orientation of the YBCO phase becomes observable which represents that these small amounts of dopant deteriorate the textured structure of the film. Also, reduction in the (008) YBCO peak intensity with increasing BZO amount is evidence for this structural deterioration. A similar effect was found by J Hänisch et al., (2006) with nanoscale precipitates of Hf dopant in YBCO films.

Figures 6, 7, 8 and 9 depict surface morphologies of undoped and doped samples. These images show that all films have a smooth, crack free surface and all of them are generally formed by c-axis oriented grains. BMeO doped YBCO films present a denser surface structure with decreasing porosity compared with the undoped YBCO films. On the other hand, 18 mol% BZO doped sample surface has bigger sized grains in comparison to the fine grains of 6 and 12 mol% BZO doped sample surfaces.

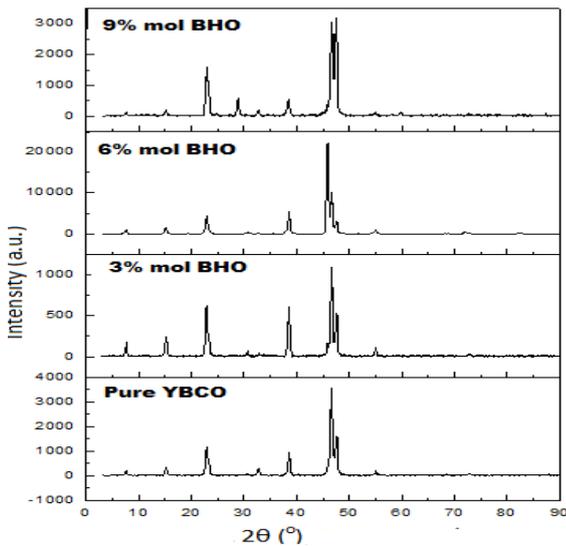


**Figure 2.** X-ray  $\theta$ - $2\theta$  scans of undoped and BZO doped YBCO films with a Co K $\alpha$  radiation.

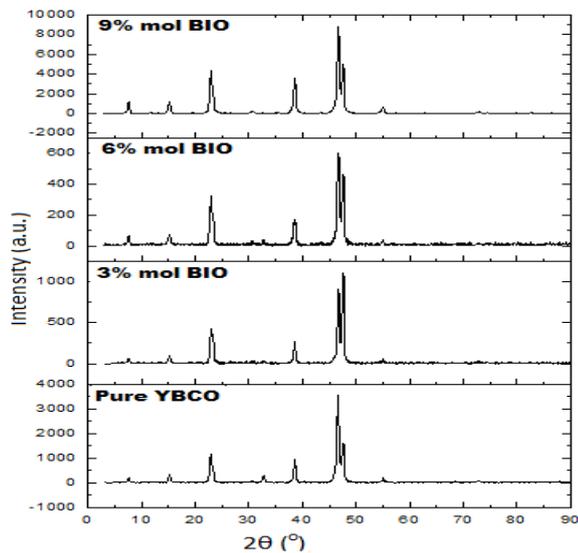
The dependence of inductively measured critical transition temperature ( $T_c$ ) on the amount of BMeO in the structure is shown in Figures 10, 11,



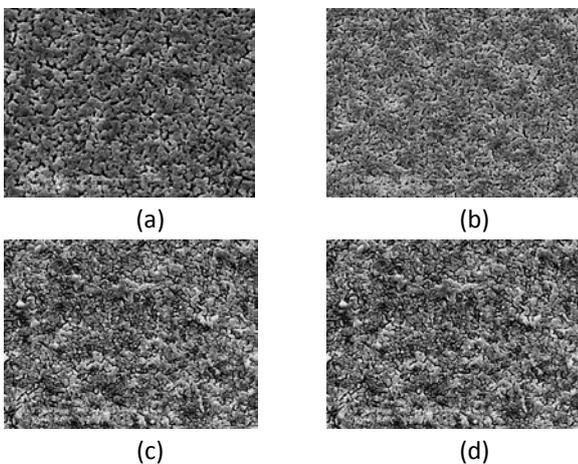
**Figure 3.** X-ray  $\theta$ - $2\theta$  scans of undoped and BMnO doped YBCO films with a Co K $\alpha$  radiation.



**Figure 4.** X-ray  $\theta$ - $2\theta$  scans of undoped and BHfO doped YBCO films with a Cu K $\alpha$  radiation.



**Figure 5.** X-ray  $\theta$ - $2\theta$  scans of undoped and BIrO doped YBCO films with a Cu K $\alpha$  radiation.

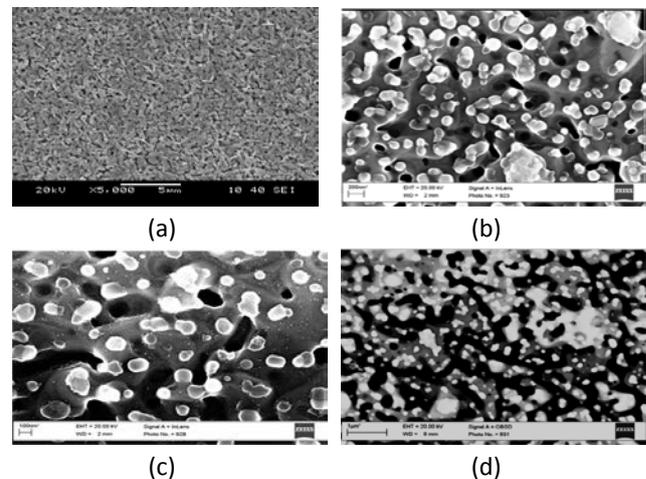


**Figure 6.** Surface morphologies of (a) pure YBCO film and YBCO films with (b) 6 mol%, (c) 12 mol% and (d) 18 mol% BZO addition.

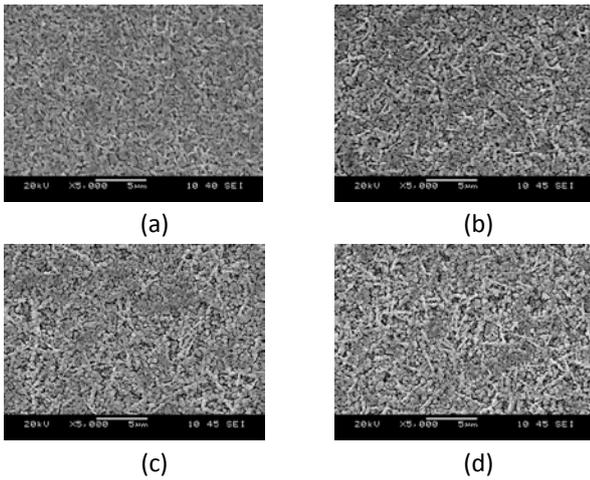
12 and 13. The  $T_c$  value of 90.1 K for an undoped YBCO film decreases gradually to 88.2 K with the increasing amount of dopant inside for BaZrO<sub>3</sub>doped films. Zr diffusion into YBCO and substitution on yttrium (Y) sites might be the reason for lowered  $T_c$ . Besides, the transition width of  $T_c$  ( $\Delta T_c$ ) stays below 1.0 K for undoped, 6 mol% and 12 mol% doped samples but is 2.0 K for 18 mol% doped sample which indicates the higher inhomogeneity of this sample.

As reported by Campbell & Evetts, (1972) the most accurate quantitative measurements of flux pinning are those of  $J_c$  as a function of field and temperature. Therefore, transport measurements were performed up to 6 T at 77 K in order to determine the magnetic field dependence of the critical current density as shown in Figure 14.

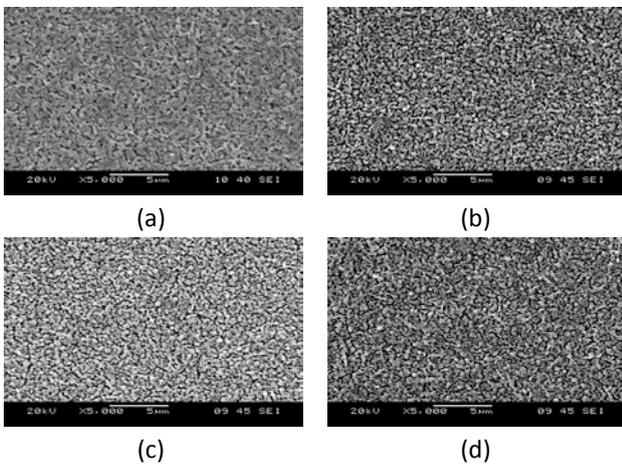
It can be seen that the 6 mol% BZO doped YBCO sample has the highest  $J_c$  value for all magnetic fields. At fields lower than 4 T, the undoped sample has higher  $J_c$  value than 12 mol% doped sample. However, the drop rate of 12 mol% doped sample is slower than the pure YBCO and  $J_c$  value of this sample exceeds the value of undoped YBCO sample at fields higher than 4 T. Within all the samples, the 18 mol% doped YBCO film has the lowest  $J_c$  value even at high fields which means that the superconducting structure of YBCO film is destroyed; self-field and in-field properties are suppressed with the excess amount of dopants concentration.



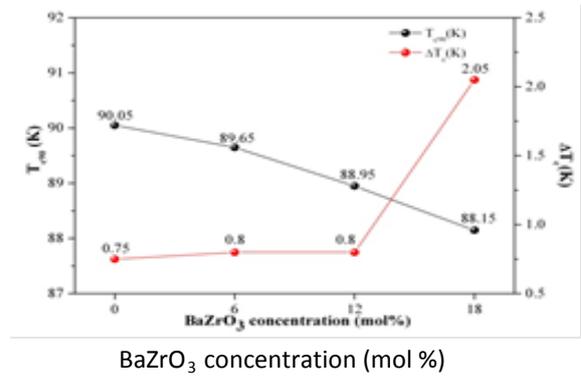
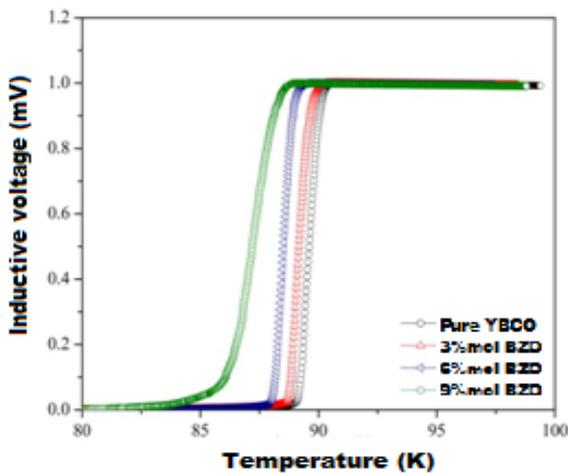
**Figure 7.** Surface morphologies of (a) pure YBCO film and YBCO films with (b) 6 mol%, (c) 12 mol% and (d) 18 mol% BMnO addition.



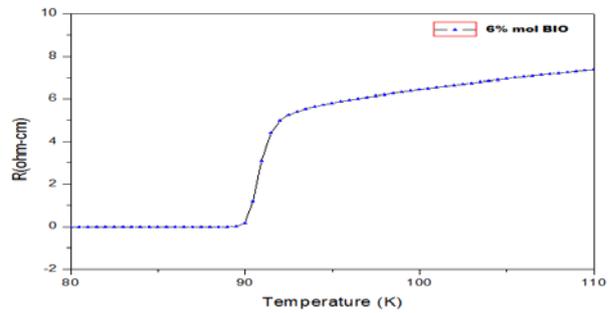
**Figure 8.** Surface morphologies of (a) pure YBCO film and YBCO films with (b) 6 mol%, (c) 12 mol% and (d) 18 mol% BHfO addition.



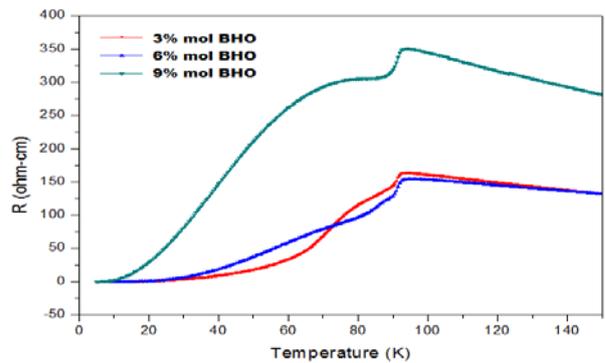
**Figure 9.** Surface morphologies of (a) pure YBCO film and YBCO films with (b) 6 mol%, (c) 12 mol% and (d) 18 mol% BIrO addition.



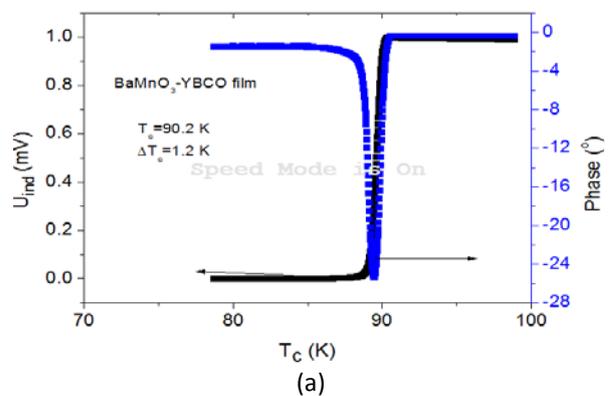
**Figure 10.** Dependence of critical temperature  $T_c$  and transition width  $\Delta T_c$  on the amount of BZO in YBCO.  $T_c$  value is decreasing with increasing amount of BZO.

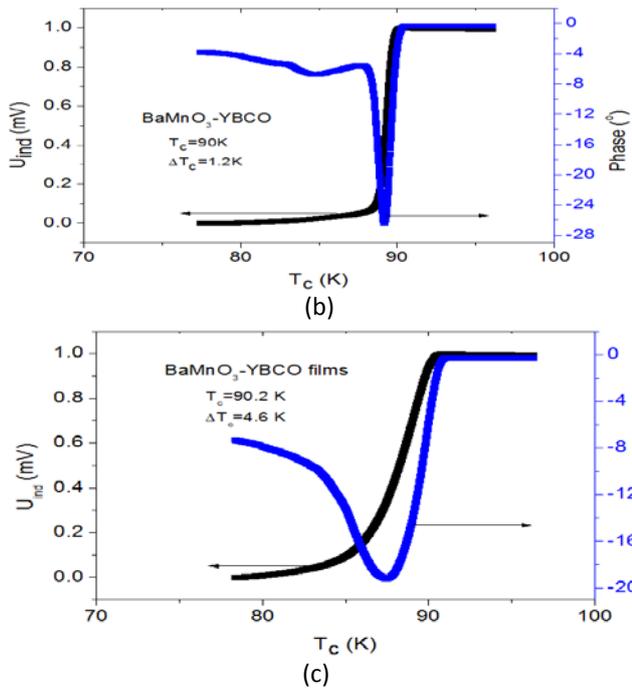


**Figure 11.** Critical temperature  $T_c$  of 6% mol BIO doped YBCO film.

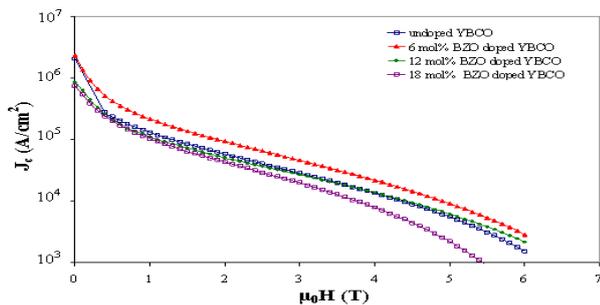


**Figure 12.** Dependence of critical temperature  $T_c$  on the amount of BHO in YBCO.





**Figure 13.** Dependence of critical temperature  $T_c$  and transition width  $\Delta T_c$  on the amount of BaMnO in YBCO.  $T_c$  value is decreasing with increasing amount of BaMnO.



**Figure 14.** Field dependence of the critical current density for all samples on a single crystal STO substrate.

#### 4. Conclusion

Barium zirconate (BZO) nanoparticles were introduced into the YBCO structure by the chemical solution deposition (CSD) method. The textured structure of undoped YBCO thin film was deteriorated with increasing dopant concentration. Both critical transition temperature ( $T_c$ ) and critical current density ( $J_c$ ) decreased with increasing dopant concentration.  $J_c$  measurements performed at 77 K to 6 T and highest  $J_c$  value has been reported for 6 mol% BZO doped sample. Further efforts will focus on improving the  $J_c$  value of films under high magnetic fields as well as the application of CSD method with other dopants and

REBCO systems.

#### Acknowledgement

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